

Linear equation in cause and effect the influence function

Defn:

In the effect at x due to a unit cause concentrated at E is denoted by the function $G(x, E)$, then the effect at x due to uniform distribution of cause of intensity $c(E)$ over an elementary region $(E + E + dE)$ is given by $c(E) G(x, E) dE$

Defn:

The effect $e(x)$ at x due to a distribution $c(E)$ over the entire region R is given by the integral

$$e(x) = \int_R G(x, E) c(E) dE \quad \text{--- (1)}$$

Defn:

The function $G(x, E)$ which represent the effect at x due to a concentrated cause at E is known as the influence function of the problem

Note:

1) If the distribution of cause is prescribed ϕ + influence function is known, then

$$e(x) = \int_R G(x, E) c(E) dE \text{ permits the determination}$$

of the effect by direct integration

2) consider a linear relation of the form $e(x) = \phi(x) + \lambda e(x)$ where ϕ is the given function or zero and λ is a constant using $e(x) = \int_R G(x, E) c(E) dE$ --- (1)

$$c(x) = \phi(x) + \lambda e(x) \quad \text{--- (2)}$$

$$\text{we get } c(x) = \phi(x) + \lambda \int_R G(x, E) c(E) dE \quad \text{--- (3)}$$

(2)

Thus relation is a Fredholm integral eqn of the second kind

using (2) in (1) we get, $e(x) = \int_0^l G(x, \xi) [\phi(\xi) + \lambda e(\xi)] d\xi$

$$E(x) = \int_0^l G(x, \xi) \phi(\xi) d\xi + \lambda \int_0^l G(x, \xi) e(\xi) d\xi \quad \text{--- (4)}$$

Both cause and effect are determined either (3) or using (4)

Problem: 1

Consider a study of small deflections of a string fixed at a point $x=0$ and $x=l$. Under load, distribution of intensity $p(x)$. Find the influence function and deflection

Soln

Suppose that the string is initially so light stretched that non-uniformity of the tension, due to small deflections, can be neglected

If a unit concentrated load is applied in the y direction at an arbitrary point ξ , the string will be deflected into two linear parts with corner at the point $x=\xi$

If we denote the uniform tension required for equilibrium in the y direction to the condition $T \sin \theta_1 + T \sin \theta_2 = 1$ --- (1)

For small deflections (and we have,

$$\left\{ \begin{aligned} \sin \theta_1, \quad x \tan \theta_1 &= \frac{\delta}{\xi} \\ \sin \theta_2, \quad x \tan \theta_2 &= \frac{\delta}{l-\xi} \end{aligned} \right\} \quad \text{--- (2)}$$

maximum deflection of the string at the loaded

Sub (2) in (1) we have

$$T - \frac{\delta}{\xi} + T \frac{\delta}{l-\xi} = 1$$

Suppose that the string is rotated uniformly the x -axis, with angular velocity ω and that in addition a continuous distribution of loading $f(x)$ is imposed in the direction radially outward from the axis of revolution.

If the linear mass density of the string is denoted by $\rho(x)$, then the total effective load intensity can be written in the form

$$P(x) = \omega^2 \rho(x) y(x) + f(x)$$

Sub $P(x)$ in (6) we get

$$\begin{aligned} y(x) &= \int_0^l G(x, \xi) [\omega^2 \rho(\xi) y(\xi) + f(\xi)] d\xi \\ &= \omega^2 \int_0^l G(x, \xi) \rho(\xi) y(\xi) d\xi + \int_0^l G(x, \xi) f(\xi) d\xi \end{aligned}$$

Problem: 2

Prove that $\omega^2 \int_0^l G(x, \xi) \rho(\xi) y(\xi) d\xi + \int_0^l G(x, \xi) f(\xi) d\xi$ is the differential equation with boundary condition $\frac{dy}{dx} \Big|_{x=0} = 0$ & $y(l) = 0$

Soln

$$\text{Given } y(x) = \omega^2 \int_0^l G(x, \xi) \rho(\xi) y(\xi) d\xi + \int_0^l G(x, \xi) f(\xi) d\xi$$

$$\text{where } G(x, \xi) = \begin{cases} \frac{x(l-\xi)}{l} & \text{if } x < \xi \\ \frac{\xi(l-x)}{l} & \text{if } x > \xi \end{cases}$$

$$\begin{aligned} y(x) &= \omega^2 \int_0^x G(x, \xi) \rho(\xi) y(\xi) d\xi + \omega^2 \int_x^l G(x, \xi) \rho(\xi) y(\xi) d\xi \\ &\quad + \int_0^x G(x, \xi) f(\xi) d\xi + \int_x^l G(x, \xi) f(\xi) d\xi \end{aligned}$$

$$y = \omega^2 \left[\int_0^x \frac{\xi(l-x)}{l} \rho(\xi) y(\xi) d\xi + \int_x^l \frac{x(l-\xi)}{l} \rho(\xi) d\xi y(\xi) \right]$$

(A)

$$\tau \delta \left(\frac{1}{\epsilon} + \frac{1}{1-\epsilon} \right) = 1$$

$$\tau \delta \left(\frac{1-\epsilon + \epsilon}{2(1-\epsilon)} \right) = 1$$

$$\tau \delta \left(\frac{1}{2(1-\epsilon)} \right) = 1$$

$$\delta = \frac{\epsilon(1-\epsilon)}{\tau \cdot 2} \quad \text{--- (3)}$$

when $x < \epsilon$ the corresponding points are $(0,0)$ & (ϵ, δ)

then the eqn is $\frac{x-0}{\epsilon-0} = \frac{y-0}{\delta-0} \quad \left[\because \frac{x-x_1}{x_2-x_1} = \frac{y-y_1}{y_2-y_1} \right]$

$$\frac{x}{\epsilon} = \frac{y}{\delta} \Rightarrow y = \frac{x\delta}{\epsilon} \quad \text{if } x < \epsilon$$

when $x > \epsilon$ the corresponding points are (ϵ, δ) & $(1,0)$

Then the eqn is $\frac{x-\epsilon}{\epsilon-\epsilon} = \frac{y-\delta}{\delta-0}$

$$\Rightarrow -\frac{y}{\delta} = \frac{x-\epsilon}{\epsilon-\epsilon} = \frac{-(\epsilon-x)}{-(\epsilon-\epsilon)}$$

$$y = \delta \left(\frac{1-x}{1-\epsilon} \right)$$

Hence the eqn of the depletion curve is

$$y = \begin{cases} \frac{x\delta}{\epsilon} & \text{if } x < \epsilon \\ \delta \frac{(1-x)}{(1-\epsilon)} & \text{if } x > \epsilon \end{cases} \quad \text{--- (4)}$$

Sub the values of δ in (4) we get

$$G(x, \epsilon) = \begin{cases} \frac{\epsilon(1-\epsilon)x}{\epsilon \cdot \tau \cdot 2} & \text{if } x < \epsilon \\ \frac{\epsilon(1-\epsilon)(1-x)}{(1-\epsilon) \tau \cdot 2} & \text{if } x > \epsilon \end{cases}$$

Hence the influence function is given by

$$= \begin{cases} \frac{x(1-\epsilon)}{\tau \cdot 2} & \text{if } x < \epsilon \\ \frac{1-x}{2} & \text{if } x > \epsilon \end{cases} \quad \text{--- (5)}$$

the depletion $y(x)$ due to a loading distribution $P(x)$ is given by $y(x) = \int_0^1 G(x, \epsilon) P(\epsilon) d\epsilon$ --- (6)

(5)

$$\begin{aligned} \frac{dy}{dx} = & \omega^2 \left[\frac{(1-x)}{\Gamma_2} \left(\int_0^x p(\varepsilon) y(\varepsilon) d\varepsilon \right)' + \frac{\varepsilon}{\Gamma_2} \int_0^x p(\varepsilon) y(\varepsilon) (-1) d\varepsilon \right] \\ & + \omega^2 \left[\frac{x(1-x)}{\Gamma_2} \left[p(\varepsilon) y(\varepsilon) \right]'_x + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} p(\varepsilon) y(\varepsilon) d\varepsilon \right] \\ & + \left[\frac{(1-x)}{\Gamma_2} \left[\varepsilon f(\varepsilon) \right]'_0 + \int_0^x \frac{\varepsilon f(\varepsilon)}{\Gamma_2} (-1) d\varepsilon \right] + \left[\frac{x(1-x)}{\Gamma_2} \left[f(\varepsilon) \right]'_x \right. \\ & \left. + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} f(\varepsilon) d\varepsilon \right] \end{aligned}$$

$$\begin{aligned} = & \omega^2 \left[\int_0^x -\frac{\varepsilon}{\Gamma_2} p(\varepsilon) y(\varepsilon) d\varepsilon + x p(x) y(x) \frac{(1-x)}{\Gamma_2} + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} \right. \\ & \left. - \frac{x(1-x)}{\Gamma_2} p(x) y(x) \right] + \int_0^1 -\frac{\varepsilon}{\Gamma_2} f(\varepsilon) d\varepsilon + \frac{x(1-x)}{\Gamma_2} f(x) \frac{dy}{dx} \\ & + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} f(\varepsilon) d\varepsilon - \frac{x(1-x)}{\Gamma_2} f(x) \end{aligned}$$

$$\begin{aligned} \frac{dy}{dx} = & \omega^2 \left[\int_0^x -\frac{\varepsilon}{\Gamma_2} p(\varepsilon) y(\varepsilon) d\varepsilon + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} p(\varepsilon) y(\varepsilon) d\varepsilon \right] \\ & + \int_0^x -\frac{\varepsilon}{\Gamma_2} + p(\varepsilon) d\varepsilon + \int_x^1 \frac{(1-\varepsilon)}{\Gamma_2} f(\varepsilon) d\varepsilon \end{aligned}$$

Now

$$\begin{aligned} \frac{d^2 y}{dx^2} = & -\omega^2 \left[\frac{\varepsilon}{\Gamma_2} p(\varepsilon) y(\varepsilon) \right]'_0 + \omega^2 \left[\frac{(1-\varepsilon)}{\Gamma_2} p(\varepsilon) y(\varepsilon) \right]'_x \\ & - \left[\frac{\varepsilon}{\Gamma_2} f(\varepsilon) \right]'_0 + \left[\frac{(1-\varepsilon)}{\Gamma_2} f(\varepsilon) \right]'_x \end{aligned}$$

$$\begin{aligned} = & \frac{\omega^2}{\Gamma_2} \left[x p(x) y(x) \right] - \frac{\omega^2}{\Gamma_2} \left[(1-x) p(x) y(x) - \frac{1}{\Gamma_2} [x - f(x)] \right. \\ & \left. - \frac{1}{\Gamma_2} [(1-x) f(x)] \right] \end{aligned}$$

$$\begin{aligned} = & \frac{+\omega^2}{\Gamma_2} \left[-x p(x) y(x) + (1-x) p(x) y(x) \right] + \frac{1}{\Gamma_2} \left[-x f(x) \right. \\ & \left. - (1-x) f(x) \right] \end{aligned}$$

$$\begin{aligned} = & \frac{\omega^2}{\Gamma_2} \left[-x p(x) y(x) - 1 [p(x) y(x) + x p(x) y(x)] \right. \\ & \left. + \frac{1}{\Gamma_2} [-x f(x) - 1 f(x) + x f(x)] \right] \end{aligned}$$

$$= \frac{\omega^2}{\Gamma_2} \left[-1 p(x) y(x) \right] + \frac{1}{\Gamma_2} [-1 f(x)]$$

⑥

$$\therefore \frac{d^2 y}{dx^2} = -\omega^2 y \quad \text{find } y(x) = -\frac{f(x)}{f}$$

multiply by f , $\therefore \frac{d^2 y}{dx^2} = -\omega^2 \text{ find } y(x) - f(x)$

$$\therefore \frac{d^2 y}{dx^2} + \omega^2 \text{ find } y(x) + f(x) = 0$$

$$y(0) = 0, y(\pi) = 0$$

Sec 8.6.

Fredholm eqns with Separable Kernels

Defn:

A kernel $K(x, t)$ separable if it can be expressed as the sum of finite no of terms each of the form is the product of a function of x alone and a function of t alone. Such a kernel is expressed in the following form

$$K(x, t) = \sum_{n=1}^N f_n(x) g_n(t)$$

Note:

In the above defn, there is without loss of generality if we assume that N functions f_n are linearly independent. Any polynomial in x is of this type.

Example:

The kernel $\sin(x-t)$ is separable

$$\text{Since } \sin(x+t) = \sin x \cos t + \cos x \sin t$$

Remark:

Integral eqns with separable kernel do occur frequently in practice. However, they are easily solved and furthermore, the result of their consideration leads to a integral eqn of

②.

general type

Also it is possible to apply the method to be developed in this section to the appropriate form of Fredholm eqn is which the kernel can be satisfactorily approximated by a polynomial in x & or a sequence of more general.

Defn:

The integral equation is said to be homogeneous if the function $p(x) = 0$

$$(i) y(x) = \lambda \int_a^b K(x, t) y(t) dt$$

Defn:

The value of λ for which $\Delta(\lambda) = 0$ are known as the characteristic value (or) Eigen values

Defn:

Any non-trivial solution of the homogeneous integral eqn is called a corresponding characteristic function (or) eigen valued function of the integral eqn

Note:

If k is of the constant $C_1, C_2, \dots, C_n$ can be assigned arbitrarily for a given characteristic value of λ then k linearly independent functions are obtained

Defn:

The integral eqn is said to be incompatible if there is no solution exist

Defn:

The integral eqn is said to be redundant if there are finite many solutions exist.

Problem:

Determine the solution of the Fredholm equation with separable kernels.

Solution:

Consider the Fredholm equation of the second kind with separable kernel.

Then it can be written as,

$$y(x) = F(x) + \lambda \int_a^b k(x, \xi) y(\xi) d\xi \rightarrow \textcircled{1}$$

$$= F(x) + \lambda \int_a^b \left[\sum_{n=1}^N f_n(x) g_n(\xi) \right] y(\xi) d\xi$$

$$= F(x) + \lambda \sum_{n=1}^N f_n(x) \int_a^b g_n(\xi) y(\xi) d\xi \rightarrow \textcircled{2}$$

[$\because k(x, \xi)$ is separable]

The Co-efficient of $f_1(x), f_2(x), \dots, f_N(x)$ in the above equations are constants.

Although their values are unknown.

$$\text{let } c_n = \int_a^b g_n(x) y(x) dx \quad (n=1 \text{ to } N)$$

$$\textcircled{2} \Rightarrow y(x) = F(x) + \lambda \sum_{n=1}^N f_n(x) c_n \rightarrow \textcircled{3}$$

This is the form of the required solution of the integral equation $\textcircled{1}$ and it remains only to determine the N constants $c_1, c_2, \dots, c_N$.

Now, on multiplying the equation $\textcircled{3}$ by $g_1(x)$ and integrating over the interval (a, b) we get,

$$\int_a^b g_1(x) y(x) dx = \int_a^b g_1(x) F(x) dx + \lambda \sum_{n=1}^N c_n \int_a^b f_n(x) g_1(x) dx$$

$$\Rightarrow C_1 = \int_a^b g_1(x) F(x) dx + \lambda \sum_{n=1}^N C_n \int_a^b f_n(x) g_1(x) dx.$$

Now, multiplying the equation (3) by $g_0(x)$ and integrating over the interval (a, b) we have.

$$\int_a^b g_0(x) y(x) dx = \int_a^b g_0(x) F(x) dx + \lambda \sum_{n=1}^N C_n \int_a^b f_n(x) g_0(x) dx$$

$$\Rightarrow C_0 = \int_a^b g_0(x) F(x) dx + \lambda \sum_{n=1}^N C_n \int_a^b f_n(x) g_0(x) dx.$$

Proceeding in this way, we have.

$$C_N = \int_a^b g_N(x) F(x) dx + \lambda \sum_{n=1}^N C_n \int_a^b f_n(x) g_N(x) dx$$

$$\text{Let } \alpha_{mn} = \int_a^b f_n(x) g_m(x) dx \text{ and } \beta_m = \int_a^b g_m(x) F(x) dx.$$

$$\therefore C_1 = \beta_1 + \lambda \sum_{n=1}^N C_n \alpha_{1n}, \quad C_0 = \beta_0 + \lambda \sum_{n=1}^N C_n \alpha_{0n},$$

$$\text{and so on. } C_N = \beta_N + \lambda \sum_{n=1}^N C_n \alpha_{Nn}.$$

$$\Rightarrow C_1 = \beta_1 + \lambda C_1 \alpha_{11} + \lambda C_0 \alpha_{10} + \dots + \lambda C_N \alpha_{1N}$$

$$C_0 = \beta_0 + \lambda C_1 \alpha_{01} + \lambda C_0 \alpha_{00} + \dots + \lambda C_N \alpha_{0N}$$

$$\vdots$$

$$C_N = \beta_N + \lambda C_1 \alpha_{N1} + \lambda C_0 \alpha_{N0} + \dots + \lambda C_N \alpha_{NN}$$

$$\Rightarrow \left. \begin{aligned} (1 - \lambda \alpha_{11}) C_1 - \lambda \alpha_{10} C_0 - \dots - \lambda \alpha_{1N} C_N &= \beta_1 \\ - \lambda \alpha_{01} C_1 + (1 - \lambda \alpha_{00}) C_0 - \dots - \lambda \alpha_{0N} C_N &= \beta_0 \\ \vdots \\ - \lambda \alpha_{N1} C_1 - \lambda \alpha_{N0} C_0 - \dots - (1 - \lambda \alpha_{NN}) C_N &= \beta_N \end{aligned} \right\} \rightarrow (4)$$

$$\begin{pmatrix} (1 - \lambda \alpha_{11}) & -\lambda \alpha_{10} & \dots & -\lambda \alpha_{1N} \\ -\lambda \alpha_{01} & (1 - \lambda \alpha_{00}) & \dots & -\lambda \alpha_{0N} \\ \vdots & -\lambda \alpha_{N0} & \dots & (1 - \lambda \alpha_{NN}) \end{pmatrix} \begin{pmatrix} C_1 \\ C_0 \\ \vdots \\ C_N \end{pmatrix} = \begin{pmatrix} \beta_1 \\ \beta_0 \\ \vdots \\ \beta_N \end{pmatrix}$$

(10)

$\Rightarrow (I - \lambda A)c = \beta$ where I is the unit matrix of order N

$$A = [a_{ij}] = \begin{pmatrix} a_{11} & a_{12} & \dots & a_{1N} \\ a_{21} & a_{22} & \dots & a_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ a_{N1} & a_{N2} & \dots & a_{NN} \end{pmatrix}, \quad c = \begin{pmatrix} c_1 \\ c_2 \\ \vdots \\ c_N \end{pmatrix} \quad \& \quad \beta = \begin{pmatrix} \beta_1 \\ \beta_2 \\ \vdots \\ \beta_N \end{pmatrix}$$

Case: (i)

Suppose $F(x) = 0$ in (1).

Then the integral equation is said to be homogeneous and obviously satisfied by the trivial solution $y(x) = 0$, corresponds to the trivial solution $c_1 = c_2 = \dots = c_N = 0$, when $\beta = 0$.

This is the only solution of determinant $\Delta = |I - \lambda A| \neq 0$. If $\Delta = 0$, atleast one of the c 's can be assigned arbitrary and the remaining c 's be determined accordingly.

Thus we have infinitely many solution of the integral equation (1).

Case: (ii)

Suppose $F(x) \neq 0$

But it's orthogonal to all function $g_1(x), g_2(x), \dots, g_N(x)$

Then $\beta = 0$.

$\therefore$ By proceeding discussion again applied to the case, except for the fact that the solution (3) of the integral equation involved also the function $F(x)$.

The trivial values $c_1 = c_2 = \dots = c_N = 0$ is the solution of $y = F(x)$.

Solution corresponding to characteristic values of λ are expressed as the sum of $F(x)$ and arbitrary multiples of characteristic functions.

Suppose atleast one right hand member of (4) does not vanish a unique non-trivial solution of the integral equation (1) is the determinant $\Delta(\lambda) \neq 0$.

If $\Delta\lambda = 0$ equation (4) are either incompatible and no solution exists or they are redundant and infinitely many solution exist.

3.7. Illustrative Examples.

Problem:

Obtain the general solution of the integral equation

$$y(x) = \lambda \int_0^1 (1 - 3x\varepsilon) y(\varepsilon) d\varepsilon + F(x).$$

Solution:

$$\text{Given } y(x) = \lambda \int_0^1 (1 - 3x\varepsilon) y(\varepsilon) d\varepsilon + F(x) \rightarrow (1).$$

$$= \lambda \int_0^1 y(\varepsilon) d\varepsilon - \lambda \int_0^1 3x\varepsilon y(\varepsilon) d\varepsilon + F(x)$$

$$y(x) = \lambda c_1 - \lambda 3x c_0 + F(x) \rightarrow (2)$$

$$\text{where } c_1 = \int_0^1 y(\varepsilon) d\varepsilon \text{ and } c_0 = \int_0^1 \varepsilon y(\varepsilon) d\varepsilon$$

Multiply (2) by 1 and integrate the result w.r. to 'x' over (0,1), we get,

$$\int_0^1 y(x) dx = \lambda \int_0^1 (c_1 - 3x c_0) dx + \int_0^1 F(x) dx.$$

$$c_1 = \lambda \left[c_1 x - \frac{3x^2}{2} c_0 \right]_0^1 + \int_0^1 F(x) dx$$

$$c_1 = \lambda \left(c_1 - \frac{3c_0}{2} \right) + \int_0^1 F(x) dx \rightarrow (3).$$

(9)

Multiply (2) by x and integrate the result w.r'to x over $(0,1)$, we get $\int_0^1 x y(x) dx = \lambda \int_0^1 [c_1 x - 3x^2 c_2] dx + \int_0^1 x F(x) dx$

$$c_2 = \lambda \left[c_1 \frac{x^2}{2} - \frac{3x^3}{3} c_2 \right]_0^1 + \int_0^1 x F(x) dx$$

$$c_2 = \lambda \left(\frac{1}{2} c_1 - c_2 \right) + \int_0^1 x F(x) dx \longrightarrow (4)$$

$$(3) \Rightarrow (1 - \lambda) c_1 + \frac{3}{2} \lambda c_2 = \int_0^1 F(x) dx \longrightarrow (5)$$

$$(4) \Rightarrow \left(-\frac{1}{2} \right) \lambda c_1 + (1 - \lambda) c_2 = \int_0^1 x F(x) dx \longrightarrow (6)$$

Now, the determinants of Co-efficient is

$$\Delta(\lambda) = \begin{vmatrix} (1 - \lambda) & 3/2 \lambda \\ -1/2 \lambda & 1 - \lambda \end{vmatrix}$$

$$= (1 - \lambda)(1 - \lambda) - \left(\frac{3}{2} \lambda \right) \left(-\frac{1}{2} \lambda \right)$$

$$= 1 - \lambda - \lambda + \lambda^2 + \frac{3}{4} \lambda^2 = 1 - \lambda^2 + \frac{3}{4} \lambda^2$$

$$= \frac{1}{4} (4 - \lambda^2)$$

$$\therefore \exists \text{ a unique solution } \Delta(\lambda) = \frac{1}{4} (4 - \lambda^2)$$

$$\Rightarrow \lambda \neq \pm 2$$

$$\text{If } \Delta(\lambda) = 0, \text{ then } \lambda = \pm 2$$

When $\lambda = 2$,

$$(5) \Rightarrow -c_1 + 3c_2 = \int_0^1 F(x) dx \longrightarrow (7)$$

$$(6) \Rightarrow -c_1 + 3c_2 = \int_0^1 x F(x) dx \longrightarrow (8)$$

When $\lambda = -2$,

$$(5) \Rightarrow 3c_1 - 3c_2 = \int_0^1 F(x) dx$$

$$\Rightarrow c_1 - c_2 = \frac{1}{3} \int_0^1 F(x) dx \longrightarrow (9)$$

$$(6) \Rightarrow c_1 - c_2 = \int_0^1 x F(x) dx \longrightarrow (10)$$

(13)

$\therefore$ ⑦ and ⑧ are incompatible unless the prescribed function $F(x)$ satisfied the condition.

$$\int_0^1 F(x) dx = \int_0^1 x F(x) dx$$

$$\Rightarrow \int_0^1 (1-x) F(x) dx = 0 \rightarrow \text{⑪}$$

⑨ and ⑩ are incompatible unless $F(x)$ satisfied the condition $\frac{1}{3} \int_0^1 F(x) dx = \int_0^1 x F(x) dx$.

$$\Rightarrow \int_0^1 \left(\frac{1}{3} - x\right) f(x) dx = 0$$

$$\Rightarrow \int_0^1 (1-3x) f(x) dx = 0 \rightarrow \text{⑫}$$

when $\lambda = \pm 2$, there exist infinitely many solution for the given integral equation.

Case: (i)

Suppose $F(x) = 0$.

Then ① is homogeneous equation.

If $\lambda \neq \pm 2$, then the only solution is a trivial one

$$y(x) = 0.$$

If $\lambda = 2$ (and $F=0$), then $c_1 = 3c_0$ [$\because$ by ⑧ & ⑦]

$$\text{Then ⑧} \Rightarrow y(x) = 2(3c_0 - 3xc_0)$$

$$= 6c_0(1-x)$$

$$= A(1-x) \quad [\because \text{characteristic function corresponding to } \lambda = 2]$$

If $\lambda = -2$, then $c_1 = c_0$ (by ⑨ & ⑩)

$$\text{⑧} \Rightarrow y(x) = -2(c_1 - 3xc_1)$$

$$= -2c_1(1-3x) = B(1-3x).$$

(14)

This is a characteristic function corresponding to $\lambda = 0$.

$\therefore$ The general solution is, $y(x) = F(x) + A_1(1-x) + B_0(1-3x)$

$$\Rightarrow \lambda(C_1 - 3xC_0) + F(x) = F(x) + A_1(1-x) + B_0(1-3x) \quad (\because \text{by } \textcircled{9})$$

$$\Rightarrow -3C_0\lambda = -A_1 - 3B_0 \quad [\because \text{equating the Co-efficients of } x]$$

Equating the Constant term, $\lambda C_1 = A_1 + B_0 \rightarrow \textcircled{14}$

$$-A_1 - 3B_0 = -3C_0\lambda$$

$$A_1 + B_0 = \lambda C_1$$

$$-2B_0 = \lambda(C_1 - 3C_0)$$

$$\Rightarrow B_0 = \frac{\lambda}{2}(3C_0 - C_1)$$

Sub in $\textcircled{14}$, $A_1 + B_0 = \lambda C_1$

$$\Rightarrow A_1 + \frac{\lambda}{2}(3C_0 - C_1) = \lambda C_1$$

$$\Rightarrow A_1 + \frac{3}{2}\lambda C_0 = \lambda C_1 + \frac{\lambda C_1}{2}$$

$$\Rightarrow A_1 = \lambda C_1 + \frac{\lambda C_1}{2} - \frac{3\lambda C_0}{2}$$

$$= -\frac{3\lambda}{2}C_0 + \frac{3}{2}\lambda C_1$$

$$A_1 = \frac{3\lambda}{2}(C_1 - C_0)$$

The general solution is $y(x) = F(x) + A_1(1-x) + B_0(1-3x)$

where $A_1 = \frac{3}{2}\lambda(C_1 - C_0)$, $B_0 = \frac{\lambda}{2}(3C_0 - C_1)$

Case: (ii)

Suppose $F(x) \neq 0$

Then there exist a unique solution $\lambda \neq \pm 0$.

If $\lambda = 0$, then $\textcircled{11}$ shows that no solution exist unless $F(x)$ is orthogonal to $(1-x)$ over $(0,1)$ in which case

(16)

infinitely many solution exist.

From ① & ⑧, $c_1 = 3c_0 - \int_0^1 F(x) dx$.

$$\therefore \textcircled{2} \Rightarrow y(x) = F(x) + \textcircled{8} \left[3c_0 - \int_0^1 F(x) dx - 3xc_0 \right]$$

$$= F(x) + 6c_0(1-x) - \textcircled{8} \int_0^1 f(x) dx$$

$$= F(x) + A(1-x) - \textcircled{8} \int_0^1 F(x) dx, \text{ where } A = 6c_0 \text{ is a}$$

Constant when $\int_0^1 (1-x)F(x) dx = 0$.

If $\lambda = -\textcircled{8}$, then ⑩ shows that no solution exist unless $F(x)$ is orthogonal to $(1-3x)$ over $(0,1)$ in which case infinitely many solution exist.

From ⑨ & ⑩, $c_1 = c_0 + \frac{1}{3} \int_0^1 F(x) dx$

$$\text{Sub } c_1 \text{ in } \textcircled{8}, y(x) = F(x) - \textcircled{8} \left[c_0 + \frac{1}{3} \int_0^1 F(x) dx - 3xc_0 \right]$$

$$= F(x) + B(1-3x) - \frac{\textcircled{8}}{3} \int_0^1 F(x) dx.$$

where $B = -\textcircled{8}c_0$ is a Constant term when

$$\int_0^1 (1-3x)F(x) dx = 0.$$